# Impact of a ACDM extension on UHECR propagation

CRPropa Workshop 25<sup>th</sup> of September

IVERSITÄT VECTOR

Janning Meinert

BERGISCHE UNIVERSITÄT OSFB1491

meinert@uni-wuppertal.de

OBSERVATORY



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meinert@uni-wuppertal.de



UNIVERSITÄT VECTOR STIFTUNG SFIT 1386



BERGISCHE UNIVERSITÄT OSFB1491





### Cosmology

### Impact of a ACDM extension on UHECR propagation Astroparticlephysics

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#### **Epochs of the Universe**



#### Epochs of the Universe



<sup>1</sup>Hofmann 2013 Nature Phys. 10.1038/nphys2793

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redshift [z]

Hofmann 2013 Nature Phys. 10.1038/nphys2793

CMB Temperature [K]

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### SU(2)<sub>CMB</sub> Fit



Hahn + Hofmann 2018, MNRAS, arxiv:1810.01253

#### **Cosmological parameters**

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	Parameter	$\Lambda CDM$ (TT,TE,EE+lowE)	SU(2)
BBN?	Z <sub>re</sub>	7.68 ± 0.79	6.23 ± 0.41
	Z <sup>*</sup>	$1089.95 \pm 0.27$	1715.9 ± 0.19
	Zp	-	52.88 ± 4.06
	$H_0 [km s^{-1} Mpc^{-1}]$	67.27 ± 0.60	74.24 ± 1.46
	$\Omega_b h^2$	0.02236 ± 0.00015	0.0173 ± 0.0002
	Ω <sub>m</sub>	0.3166 ± 0.0084	0.384 ± 0.006
	σ <sub>8</sub>	0.8120 ± 0.0073	0.709 ± 0.020
	$S_8 \equiv \sigma_8 (\Omega_m / 0.3)^{0.5}$	0.834 ± 0.016	0.8021 ± 0.0227
	$\Omega_{old}$	_	0.113 ± 0.002
	$\Omega_{new}$	-	$0.0771 \pm 0.0012$

Hahn + Hofmann 2018, MNRAS, arxiv:1810.01253





## How is T(z) modiefied?

### Friedmann-Lemaître-Robertson-Walker Universe

energy density energy pressure  $\frac{\mathrm{d}\rho}{\mathrm{d}a} = -\frac{3}{a} \left( \rho + P \right) , \quad s = \frac{\rho + P}{T} .$ continuity eq.:  $a = \exp\left(-\frac{1}{3}\log\frac{s(T)}{s(T(z=0))}\right)$  $\frac{s(T)}{s(T(z=0))} = \frac{8}{2} \left(\frac{T}{T(z=0)}\right)^3, \ (T \gg T(z=0))$  $a = \frac{1}{z+1} = \left(\frac{1}{4}\right)^{\frac{1}{3}} \frac{T_0}{T} \quad \text{Changed T(z)}$ 

### How do we change $T(z) = S(z) (1+z) T_0$



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### How do we change $T(z) = S(z) (1+z) T_0$



### Isn't T(z) measured already?

How do we measure T(z)?



How do we measure T(z)?



### Summary of T(z) effect

#### **CMB** Evolution



#### **CMB** Evolution



#### **CMB** Evolution



# Propagation of UHECRs



redshift [z]





### CR spectrum 1<sup>st</sup> Knee



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CR spectrum

1<sup>st</sup> Knee



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### UHECR propagation with Prince-CR



### **UHECR** propagation









####


#### **UHECR** propagation



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# Hadronic interactions

#### Hadronic Interaction models



#### Hadronic Interaction models



# Cosmogenic neutrinos

#### **Cosmogenic Neutrinos**



#### **CMB** Evolution



#### How do we change $T(z) = S(z) (1+z) T_0$



#### Overview of different T(z)



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Overview of different T(z)
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#### Summary



JM et al. (2023) 2309.08451

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#### Thank you!











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# Extra slides

# Screening effects due to massive gauge modes









SU(2)





# S8 tension

#### S8 Tension

 CMB Planck CMB TT, TE, EE+lowE CMB Planck CMB TT, TE, EE+lowE+lensing · CMB ACT+WMAP

· WL KIDS-1000 Asgari et al. (2021) \* WL KIDS+VIKING+DES-Y1 Asgari et al. (2020) \* WL KIDS+VIKING+DES-YI Joudaki et al. (2020) \* WL KEDS+VIKING-450 Wright et al. (2020) \* WL KIDS+VIKING-450 Kohlinger et al. (2017) + WL KIDS-450 • WL KiDS-450 · WL DES-Y3 · WL DES-YI Troxel et al. (2018) · WL BSC-TPCF Hamana et al. (2020) · WL HSC-pscudo--C Hikage et al. (2019) · WL CFHTLenS Joudaki et al. (2017) \* WL+GC HSC+BOSS Miyatake et al. (2022) \* WL+GC DESI+CMB White et al. (2022) \*WL+GC+CMBL KIDS+DES+eBOSS+DEL5 \* WL+GC KiD5-1000 3x2pt Heymans et al. (2021) \* WL+GC KilD5-450 3x2pt Joudaki et al. (2018) \* WL+GC DES-Y3 3x2pt Abbott et al. (2021) \* WL+GC DES-YI 3x2pt Abbott et al. (2018d) WL+GC KiDS+VIKING-450+BOSS Tröster et al. (2020) \* WL+GC KiDS+GAMA 3x2pt van Uitert et al. (2018) · GC BOSS DR12 bispectrum Philcox et al. (2021) GC BOSS+eBOSS Ivanov et al. (2021) · GC BOSS power spectra Chen et al. (2021) \* GC BOSS DR12 Tröster et al. (2020) \* GC BOSS Galaxy Power Spectrum Ivanov et al. (2020) \* GC+CMBL unWISE+Planck CMB lensing Krolewski et al. (2021) CC AMICO KIDS-DR3 Lesci et al. (2021) · CC DES-YI Abbott et al. (2020d) CC SDSS-DR8 Costanzi et al. (2019) ·CC XMM-XXL Pacatal et al. (2018) CC ROSAT (WIG) Mantz et al. (2015) · CC SPT ISZ Bocquet et al. (2019) · CC Planck tSZ Salvati et al. (2018) · CC Planck tSZ Ade et al. (2016d) · RSD Benisty (2021) · RSD 0.2 0.4 0.6 0.8 1.0 1.2  $S_8 \equiv \sigma_8 \sqrt{\Omega_m} / 0.3$ 

Aghanim et al. (2020d) Aghanim et al. (2020d) Aiola et al. (2020)

0.834 0.832 0.84

#### Early Universe

Late Universe

Hildebrandt et al. (2020) Hildebrandt et al. (2017) Amon et al. and Secco et al. (2021) Garcia-Garcia et al. (2021)

Kazantzidis and Perivolaropoulos (201

### BBN??



Tanabashi+ PRD 2018

# Why SU(2)CMB?

#### CMB radio excess



Seiffert et al. (2011)

#### ARCADE 2



Temperature = 2.725 K



# Isn't T(z) measured already?

<sup>2</sup> To a very good approximation the spectral intensity  $I(\nu)$  of today's CMB is given as  $I_{z=0}(\nu)d\nu = 16\pi^2 \frac{\nu^3}{\exp\left(\frac{2\pi\nu}{T(z=0)}\right)-1}d\nu$  Mather et al. (1994). If we assume a T-z relation of  $T(z = 0) = \frac{1}{f(z)}T(z)$  and a  $\nu$ -z relation of  $\nu(z = 0) \equiv \frac{1}{g(z)}\nu'$  with  $f(z) \neq g(z)$ , then the Stefan-Boltzmann law would still have redshifted according to the T-z relation:  $\int d\nu I_{z=0}(\nu) = \frac{\pi^2}{15} T^4(z=0) = \frac{\pi^2}{15} \left(\frac{T(z)}{f(z)}\right)^4 = \left(\frac{1}{g(z)}\right)^4 \int d\nu' I_z(\nu').$ However, the maximum  $v_{\text{max}} = \frac{2.821}{2\pi}T(z = 0)$  of the distribution  $I_{z=0}(v)dv$  converts to a maximum  $v'_{\max} = \frac{2.821}{2\pi} \frac{g(z)}{f(z)} T(z)$  of the distribution  $I_z(\nu')d\nu' = 16\pi^2 \frac{(\nu')^3}{\exp\left(\frac{f(z)}{\sigma(z)}\frac{2\pi\nu'}{T(z)}\right) - 1}d\nu'$ . Thus,  $I_z(\nu')$  no longer

would be a blackbody spectrum.

e-Print: 2303.16744 [hep-th]

# Energy loss length

#### Changed Energy Loss Length



# Lensing

#### Cosmological parameters

Parameter	$\Lambda CDM$ (TT,TE,EE+lowE)	SU(2)
Z <sup>*</sup>	$1089.95 \pm 0.27$	$1715.9 \pm 0.19$
Ω <sub>m</sub>	$0.3166 \pm 0.0084$	$0.384 \pm 0.006$
σ <sub>8</sub>	$0.8120 \pm 0.0073$	$0.709 \pm 0.020$



KiDS Collaboration arxiv.org/pdf/2007.15633

#### $\sigma_8$ tension



https://kids.strw.leidenuniv.nl/cosmicshear2016.php
## $H_0$ tension



distance ladder





CMB Temperature [K]

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