

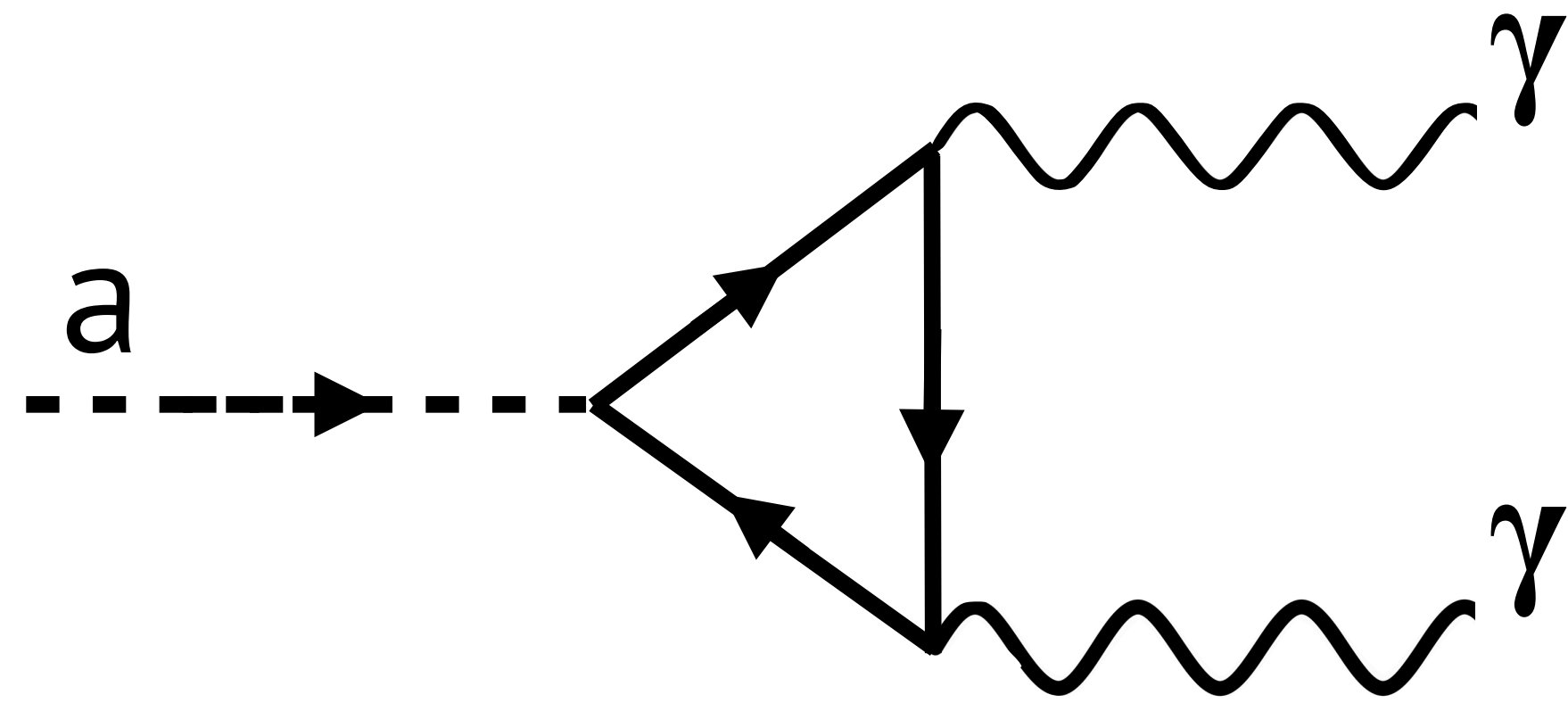
coupling axion-like particles to photons during propagation: the ALP_{inist} plug-in

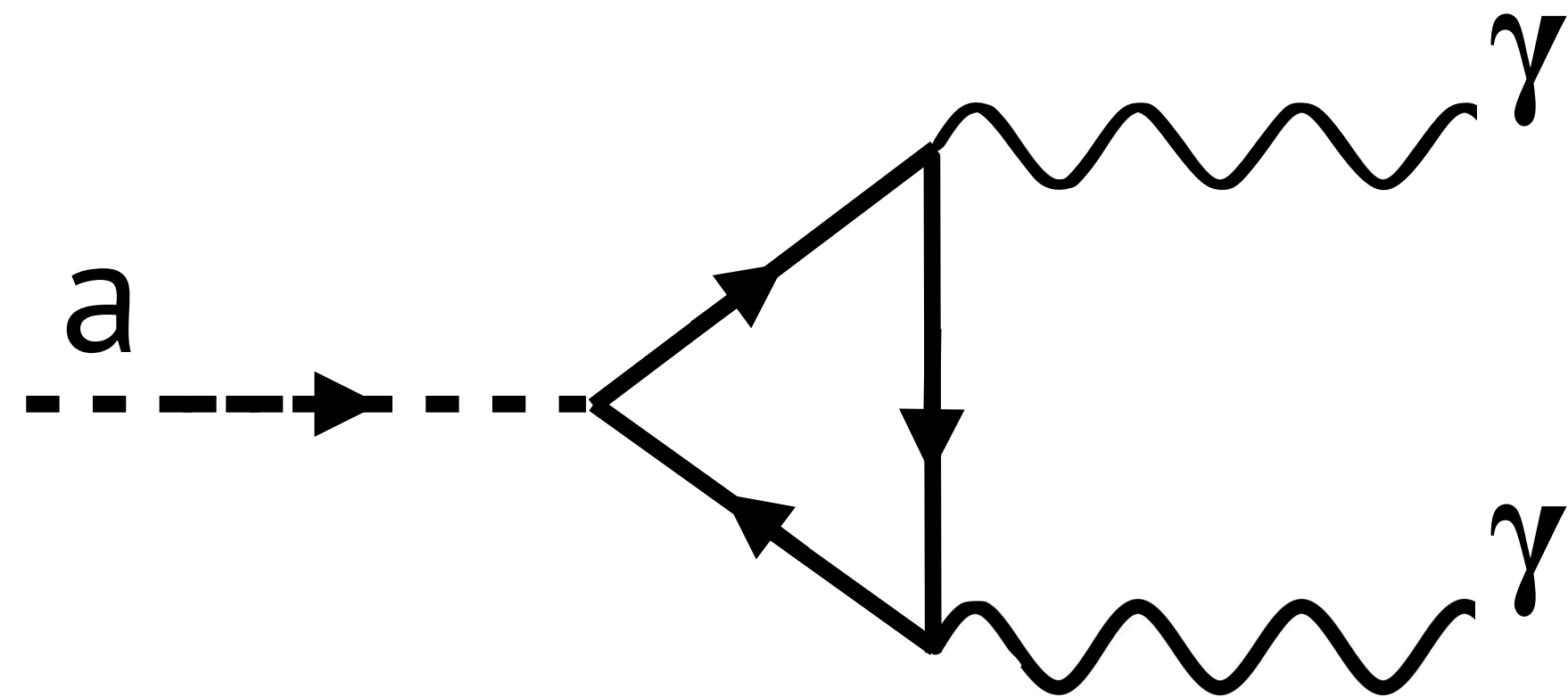
Rafael Alves Batista, Cristina Viviente, Gaetano Di Marco, Miguel A. Sánchez-Conde

Instituto de Física Teórica UAM-CSIC, Madrid, Spain

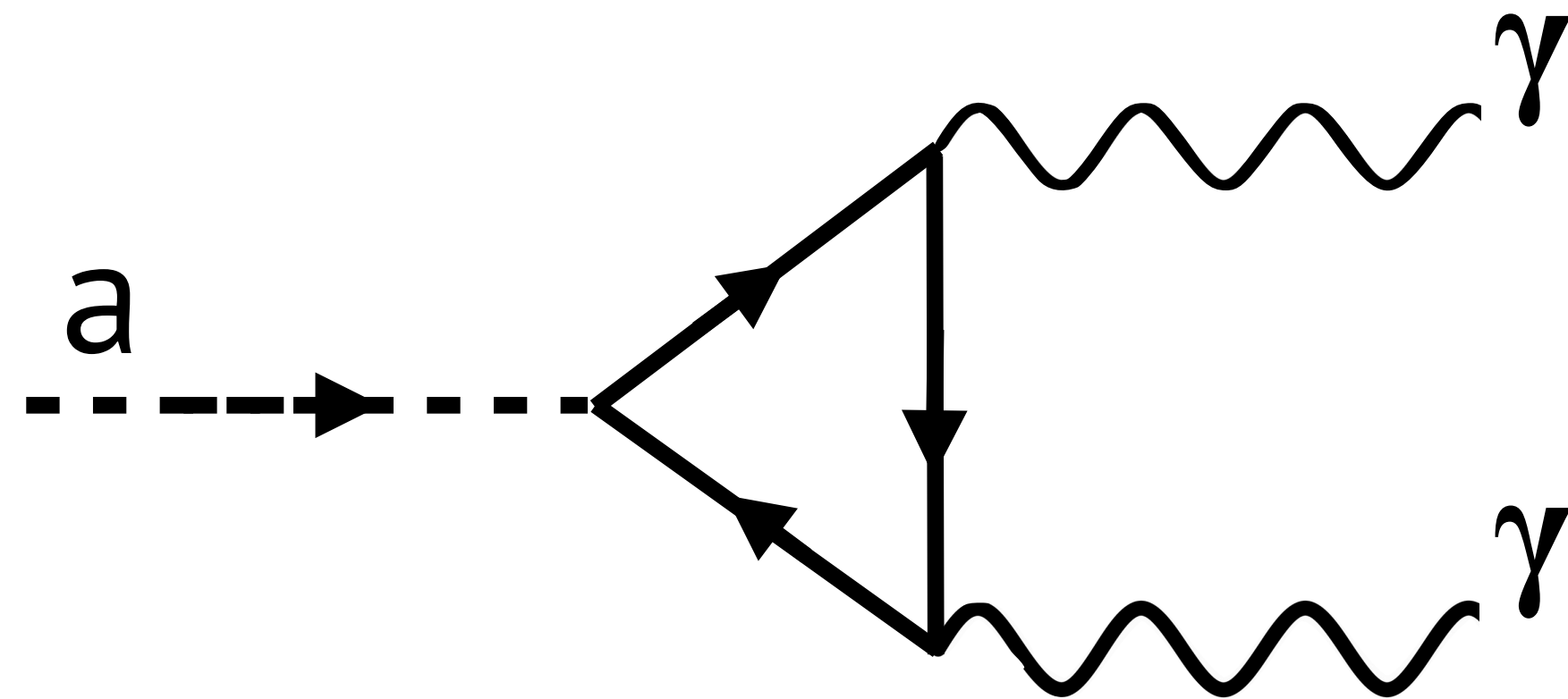
Departamento de Física Teórica, Universidad Autónoma de Madrid, Madrid, Spain

rafael.alvesbatista@uam.es



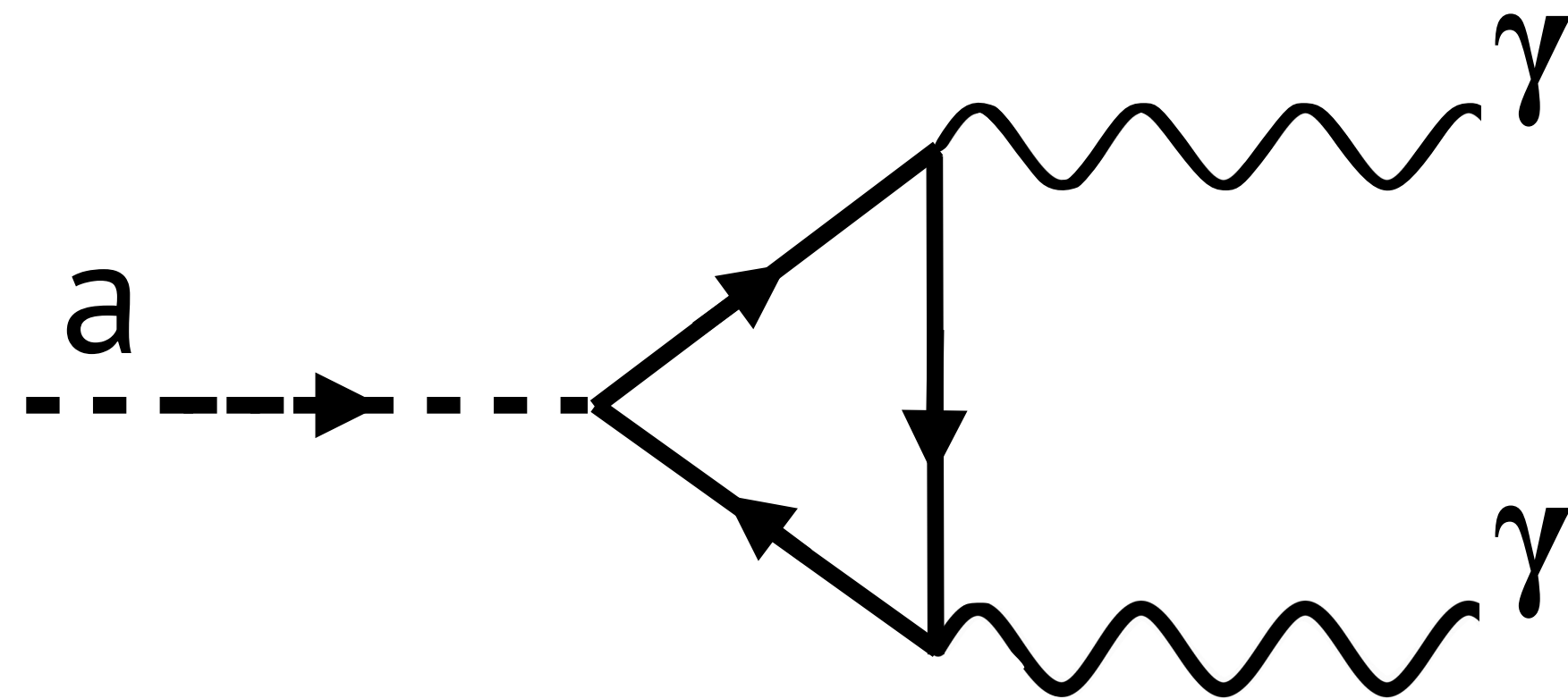


Lagrangian ALP-photon mixing



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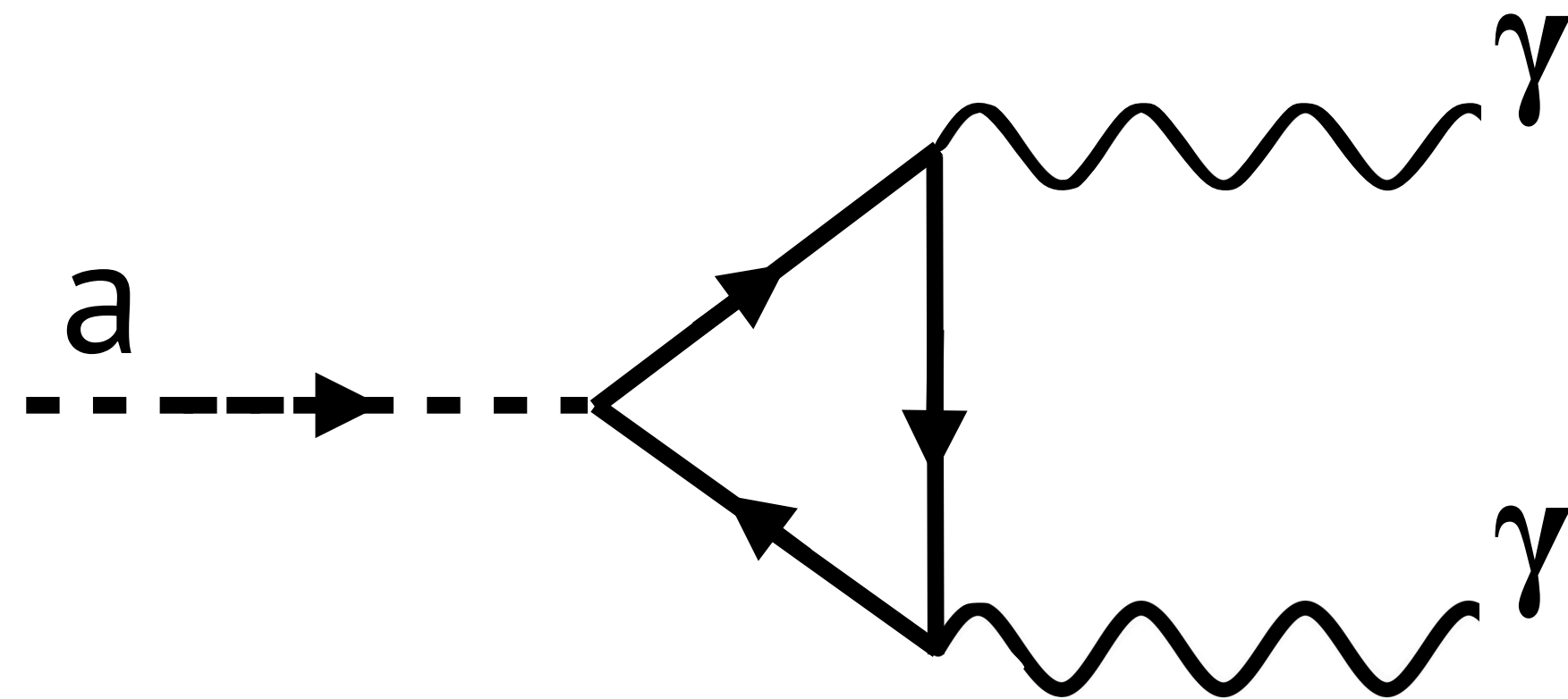
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coupling
constant

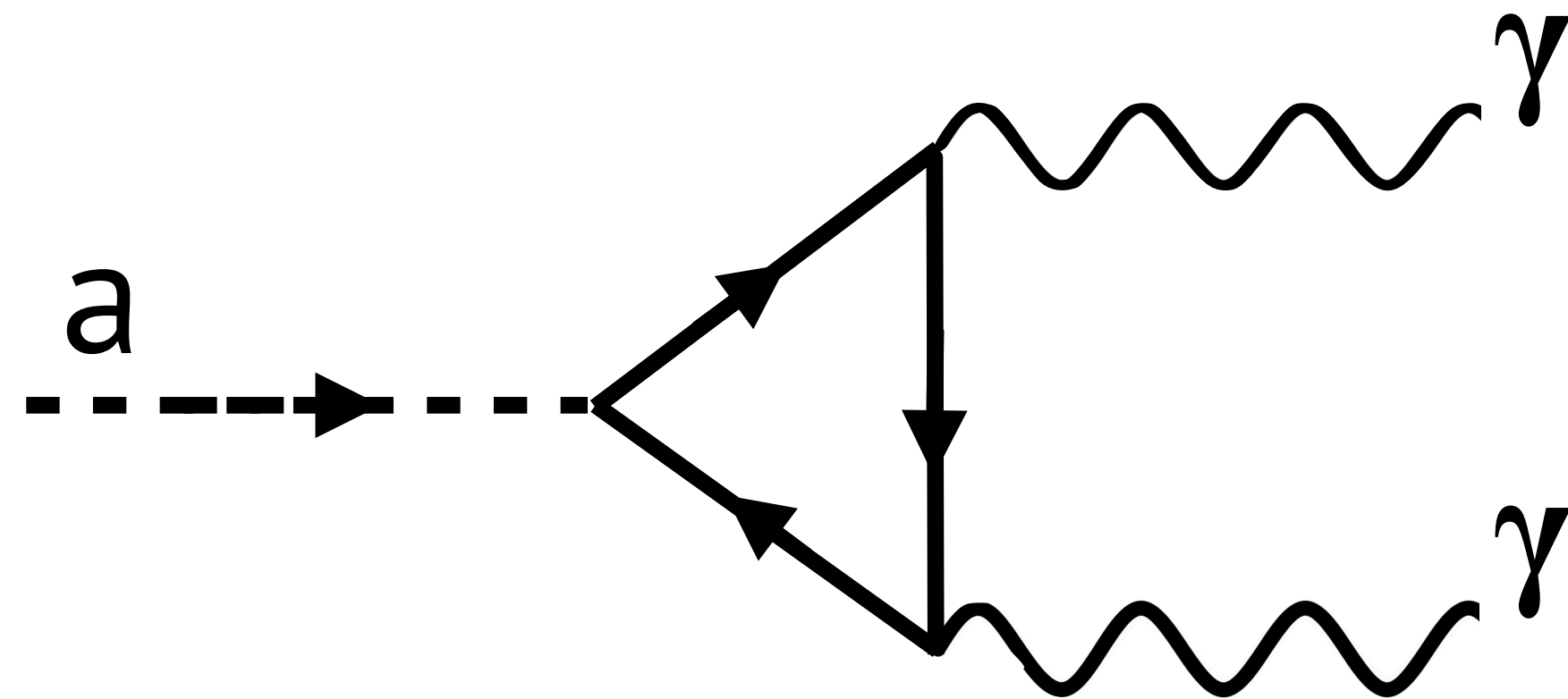


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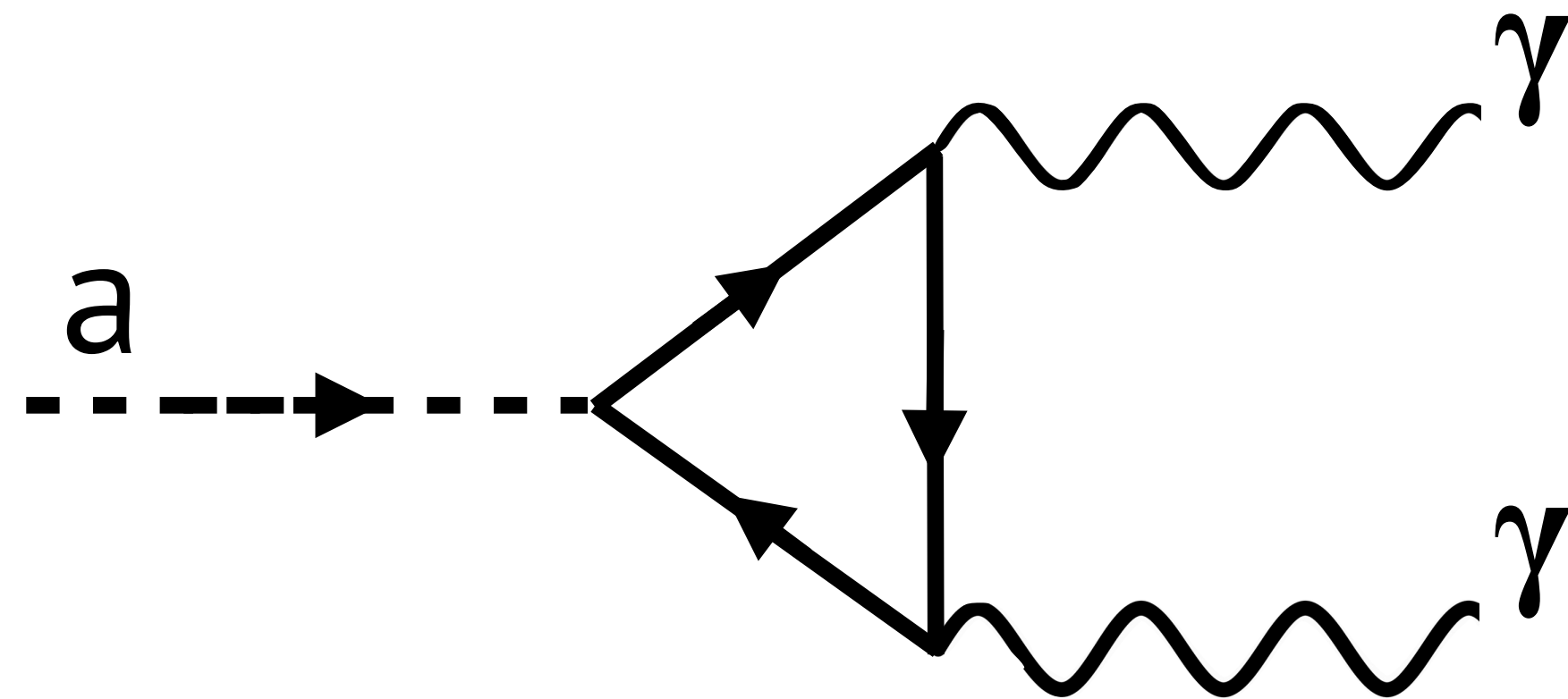
photon polarisation



Lagrangian ALP-photon mixing

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coupling constant photon polarisation external magnetic field



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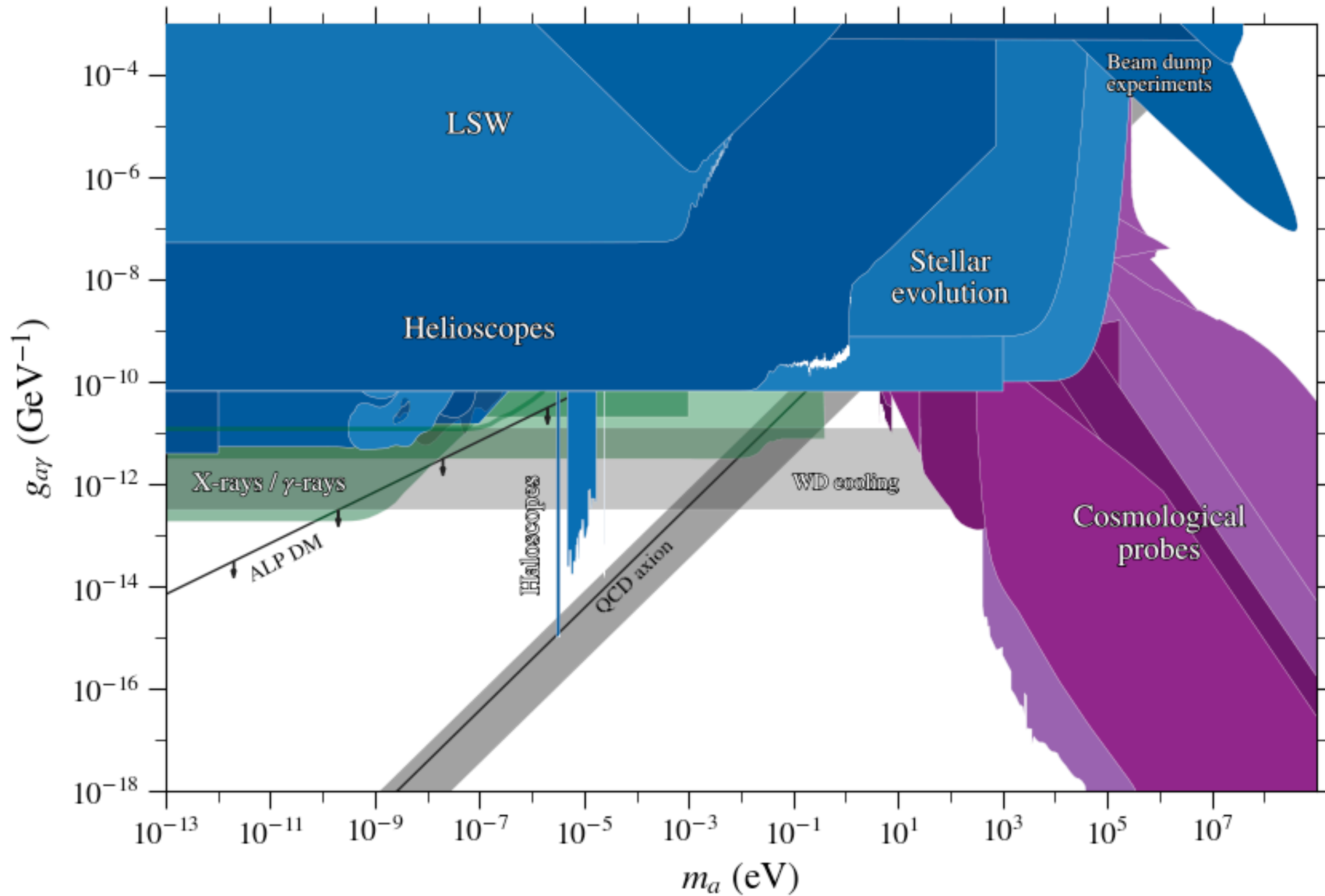
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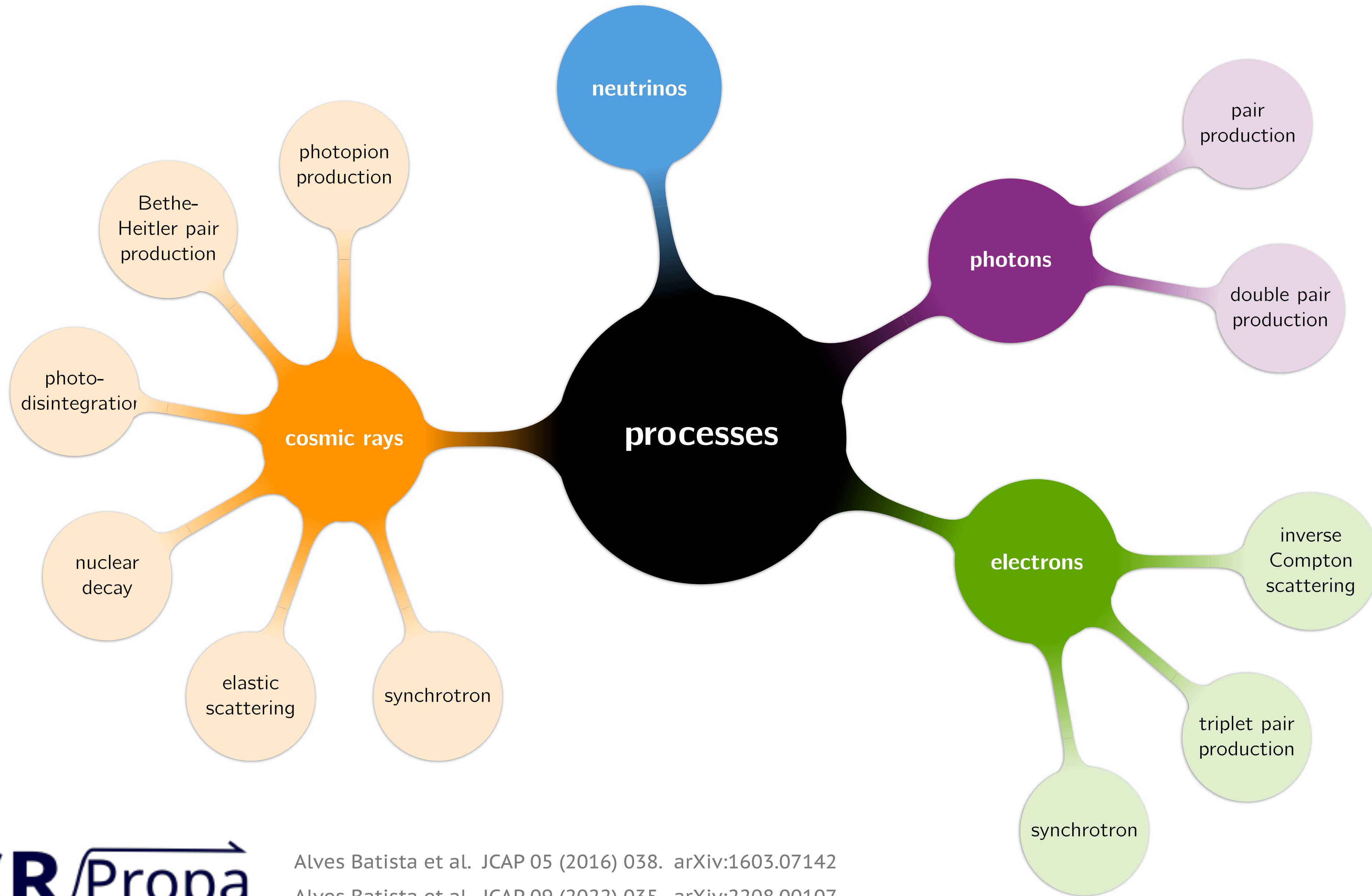
ALP field

axion-like particles: parameter space

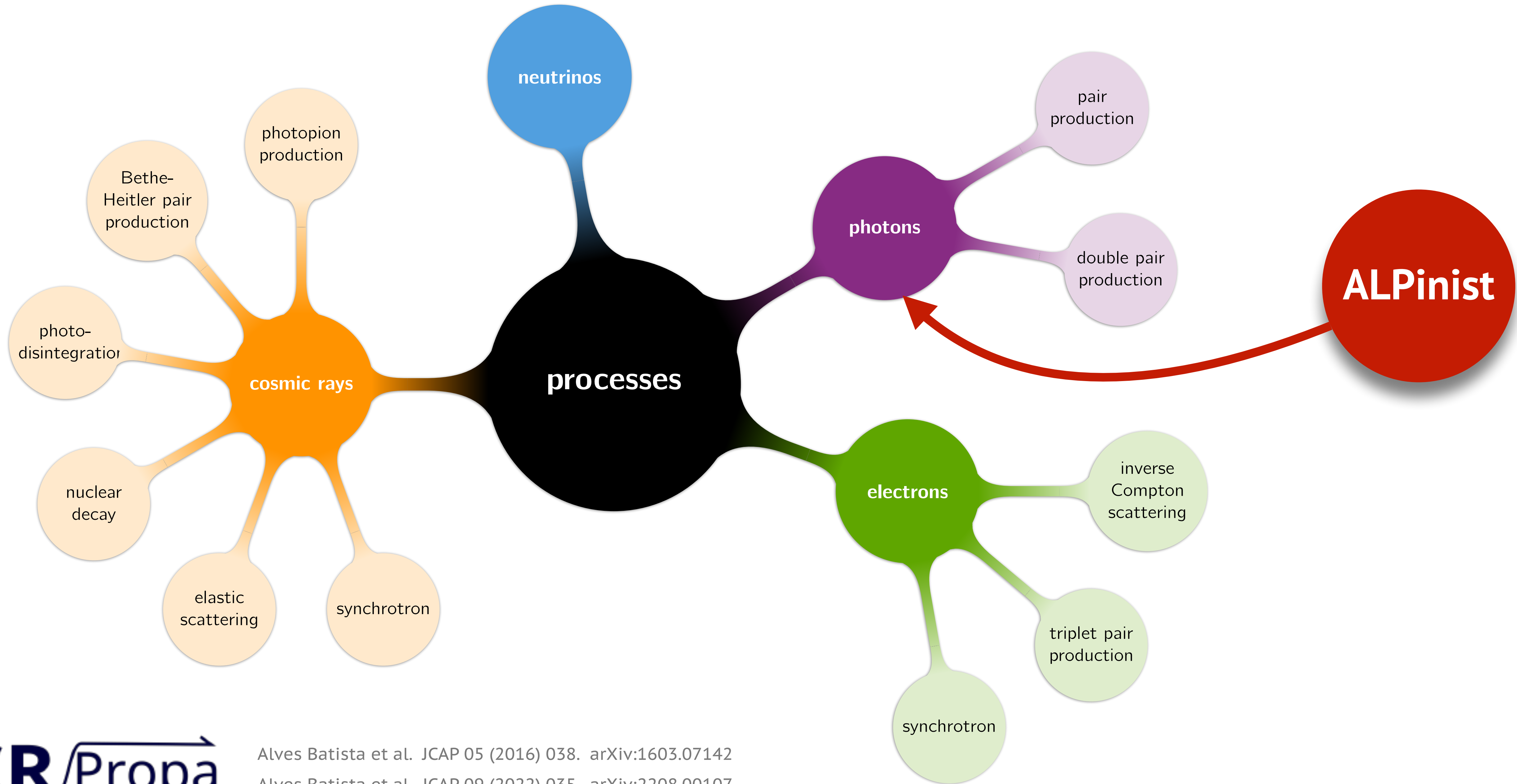


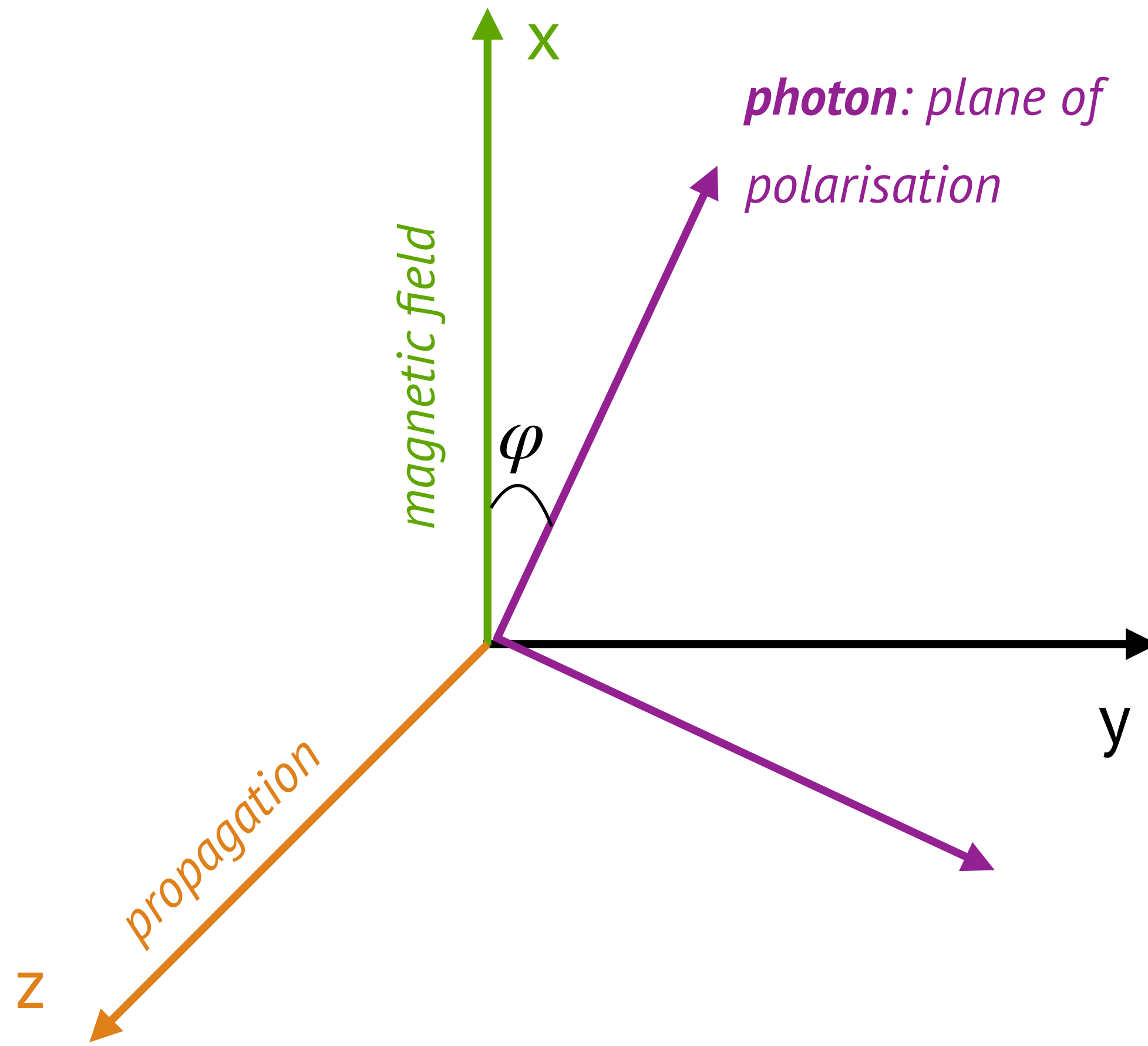
<https://github.com/me-manu/gammaALPsPlot>

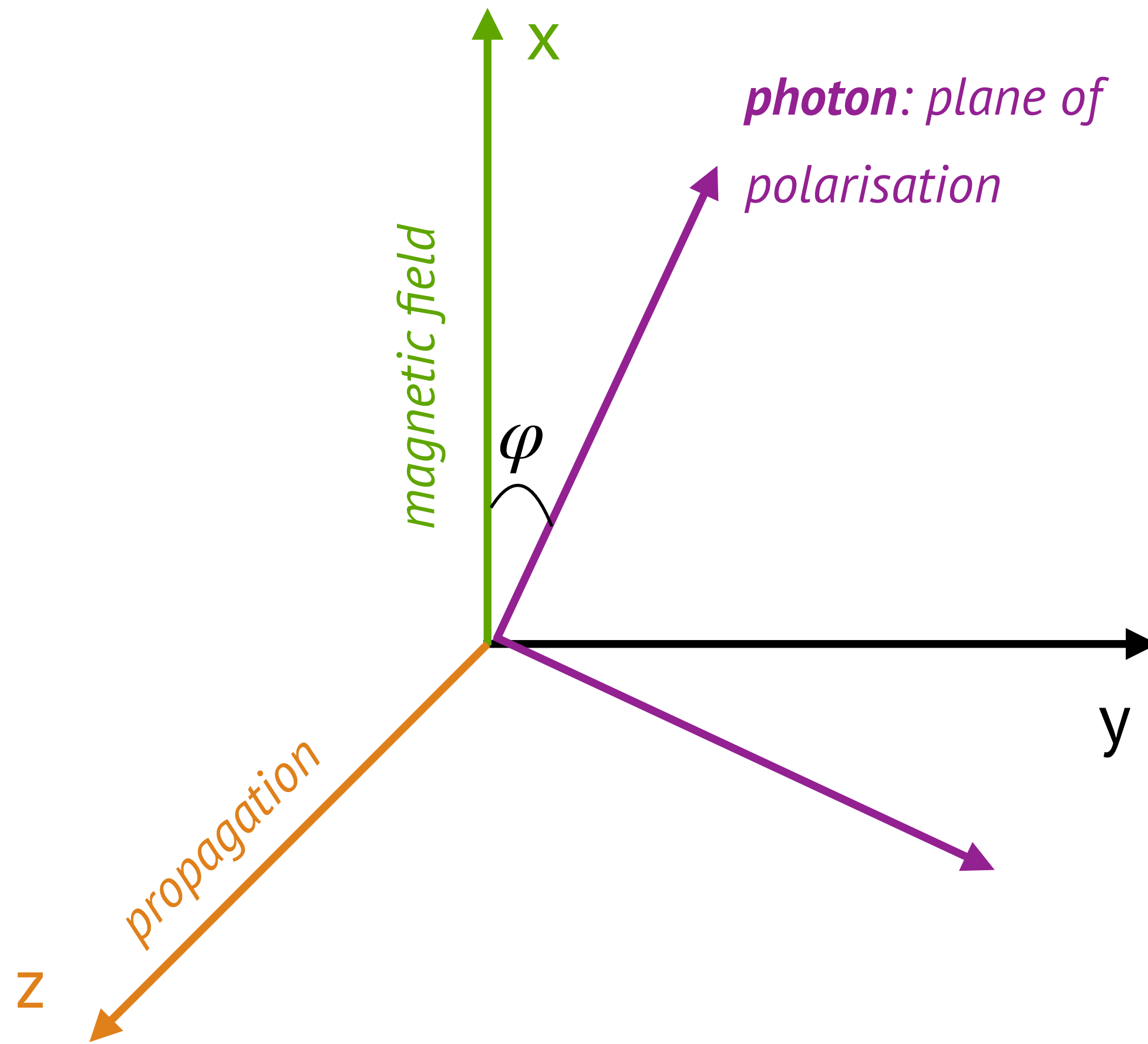
Monte Carlo simulations of ALP-photon mixing



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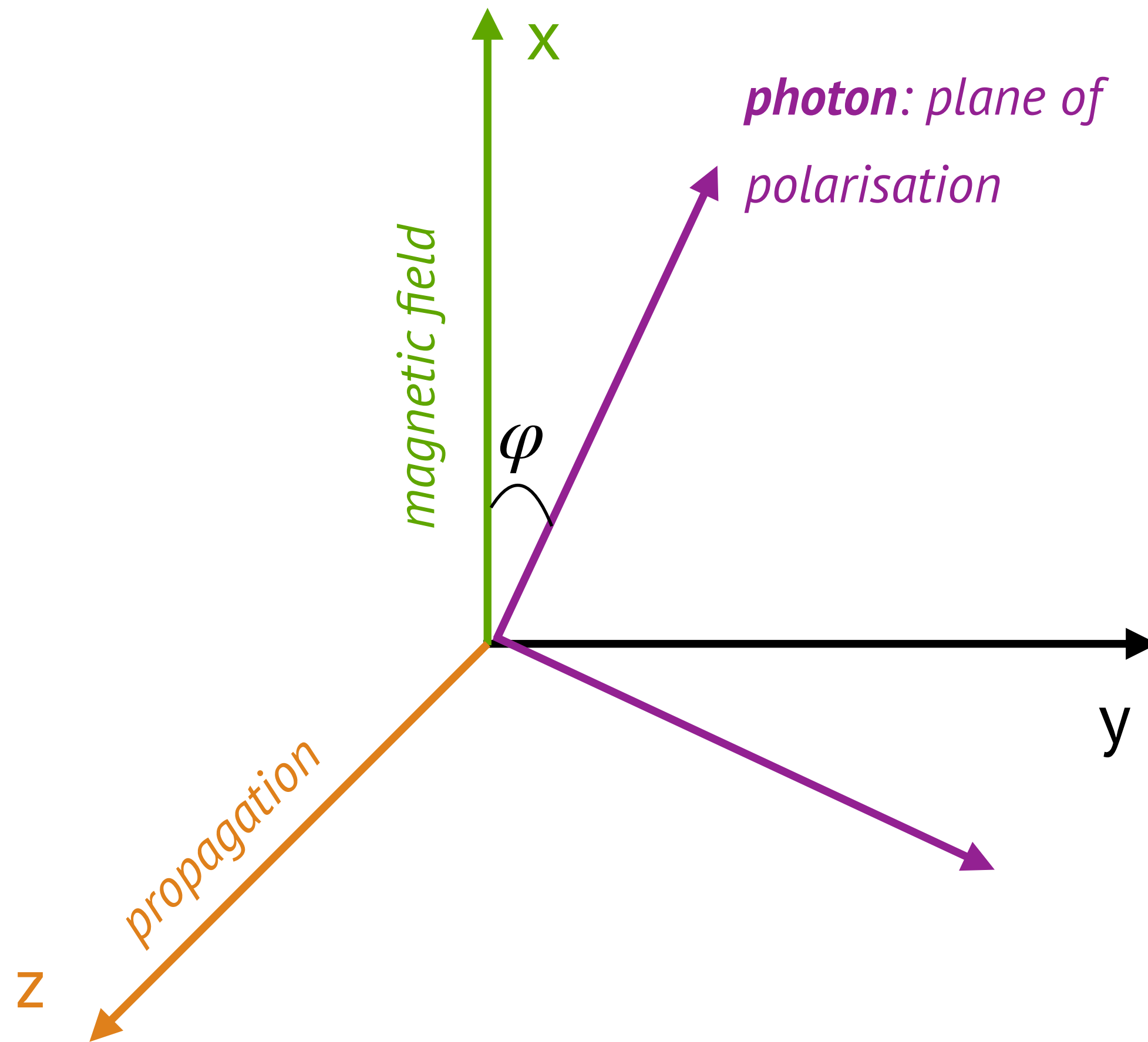






equation of motion

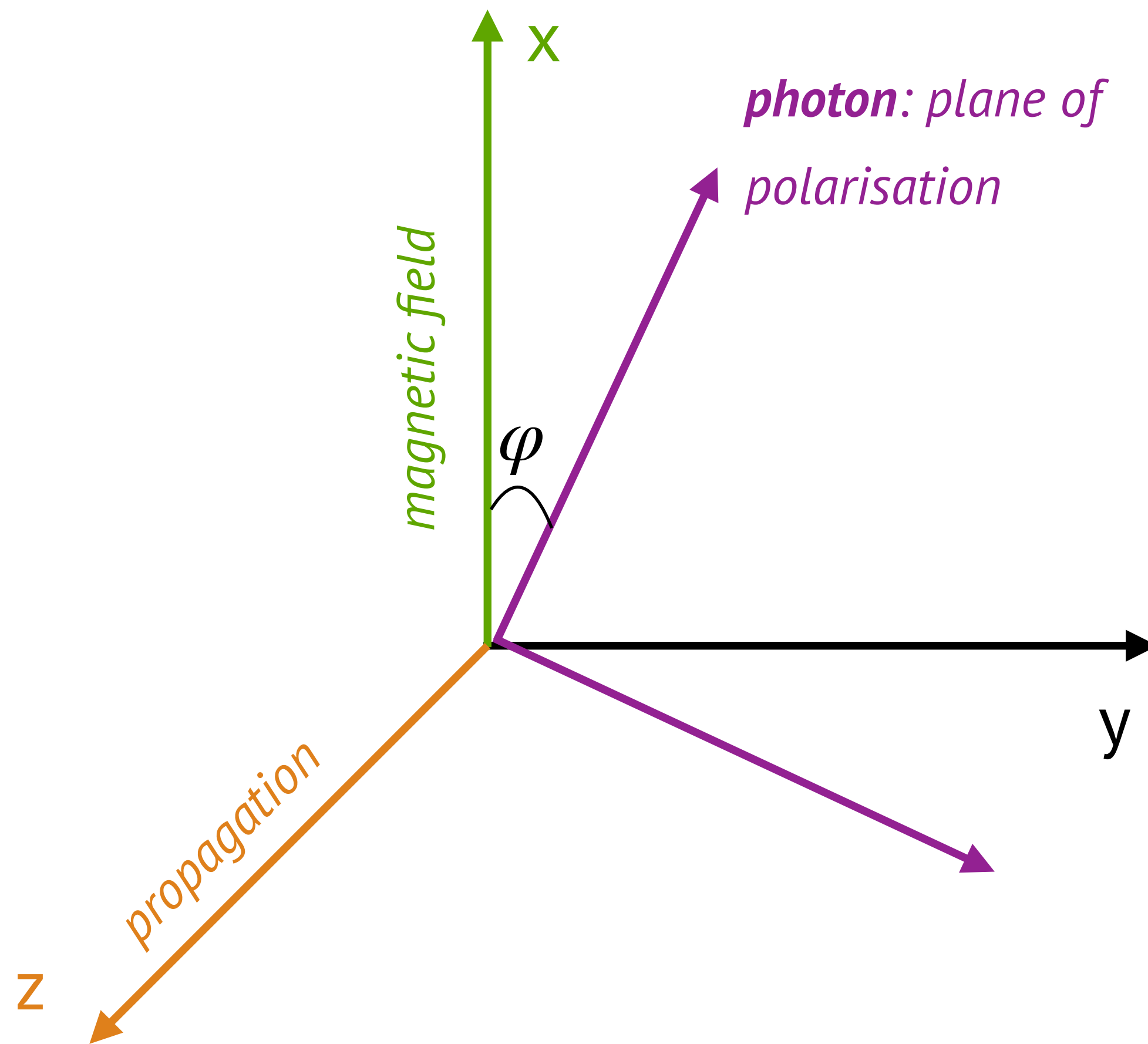
$$\left(E - i \frac{\partial}{\partial z} + \frac{1}{2E} \mathbb{M} \right) \vec{\mathcal{A}} = 0$$



three-state quantum system

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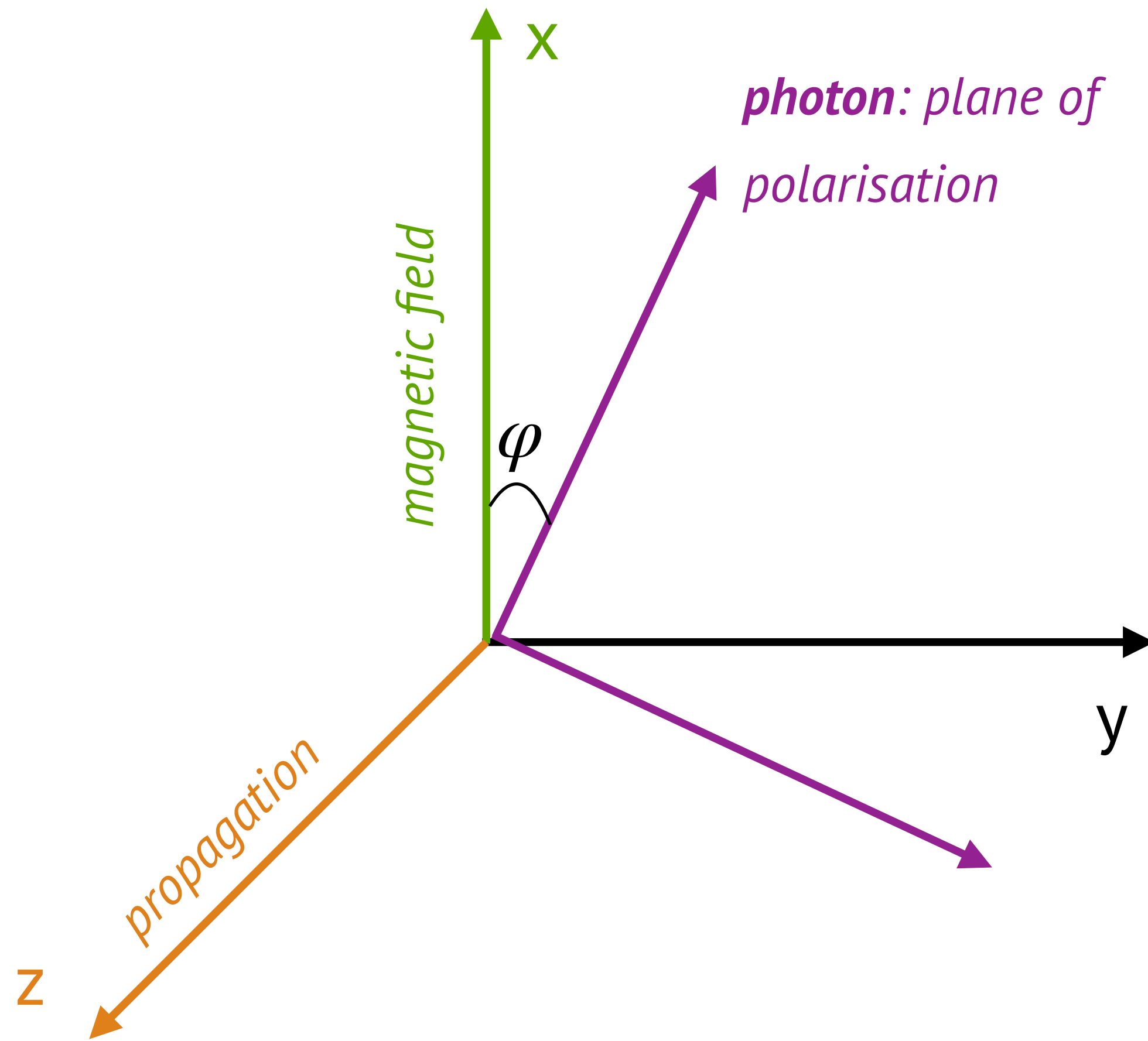


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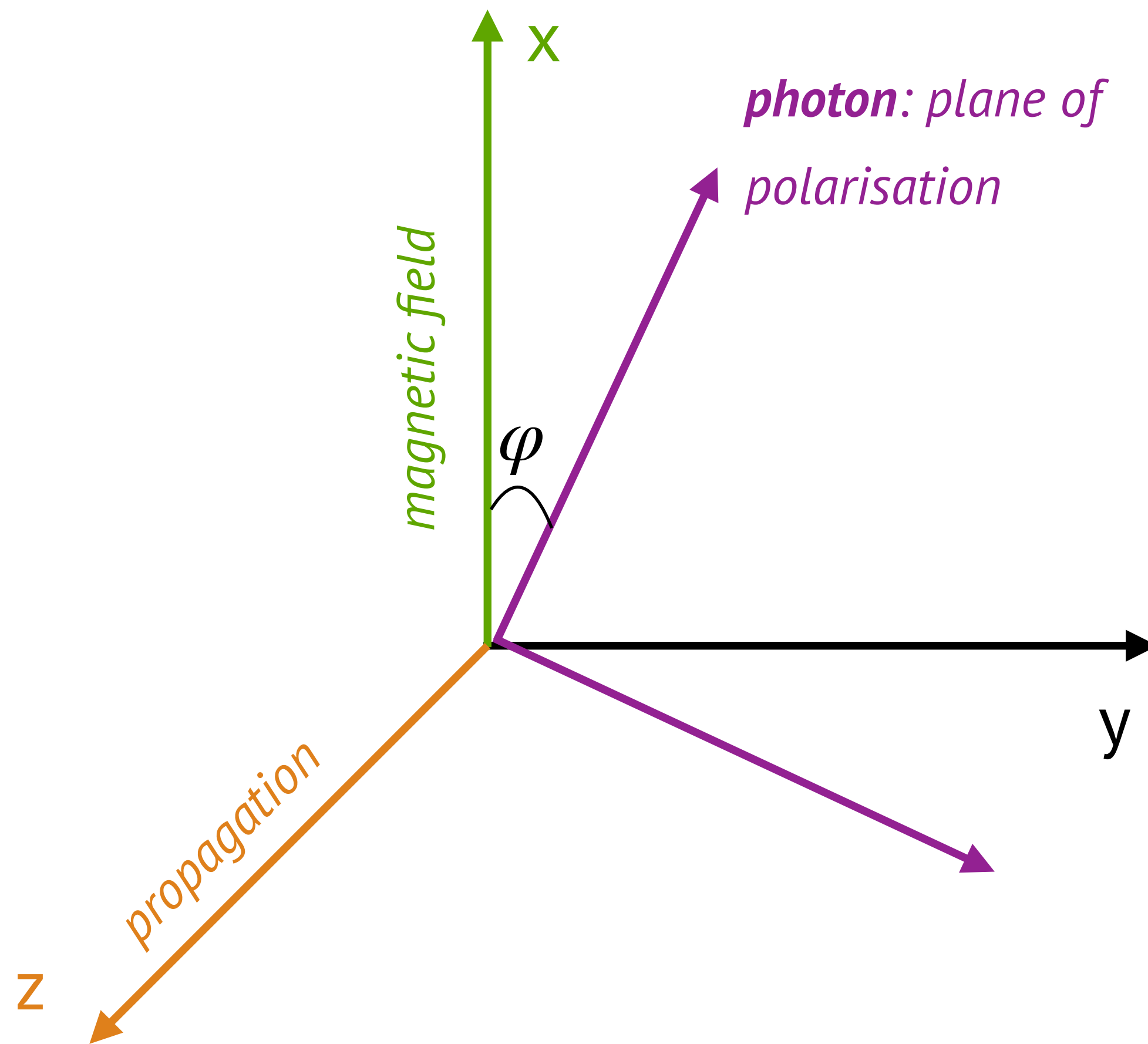


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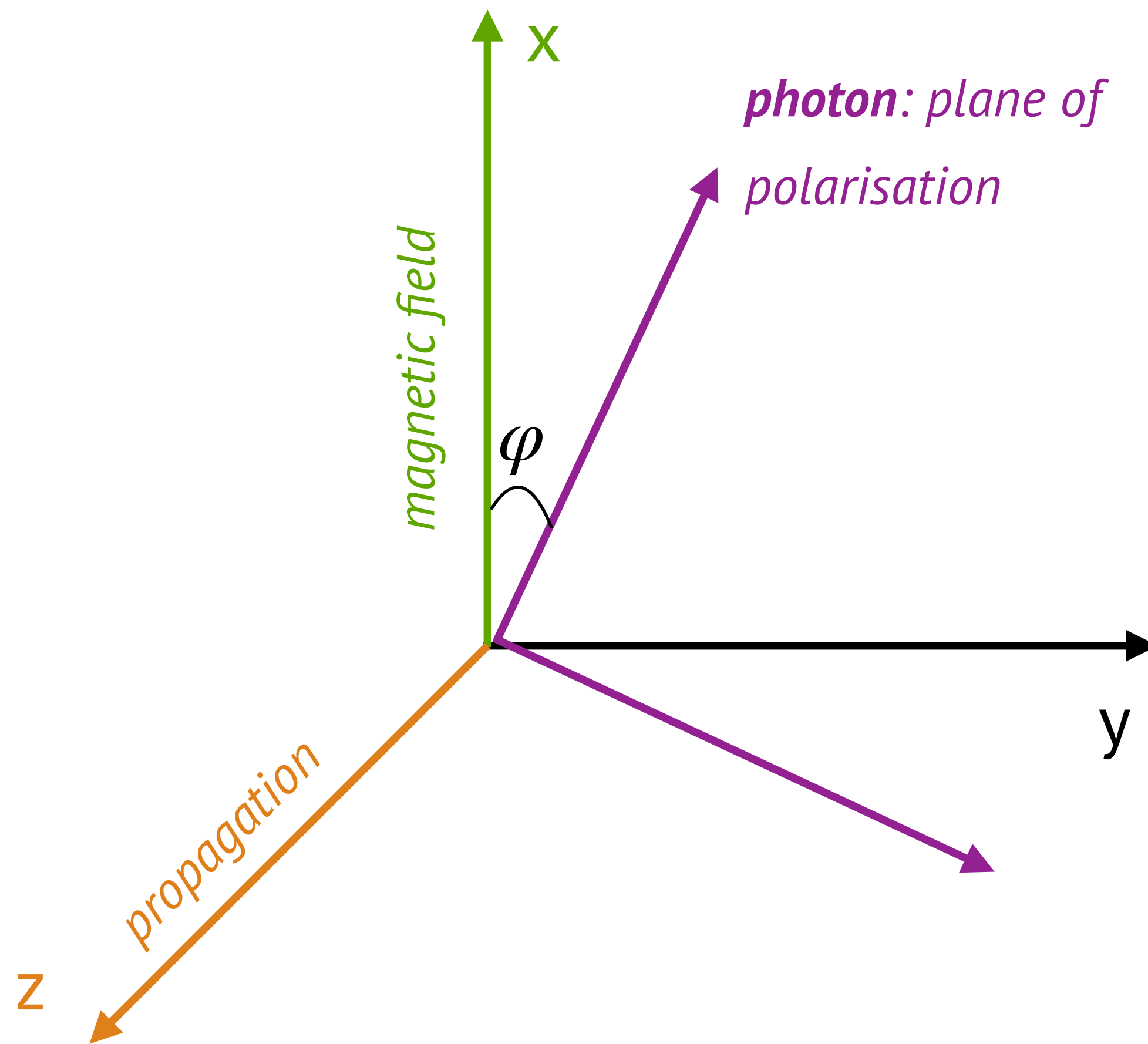


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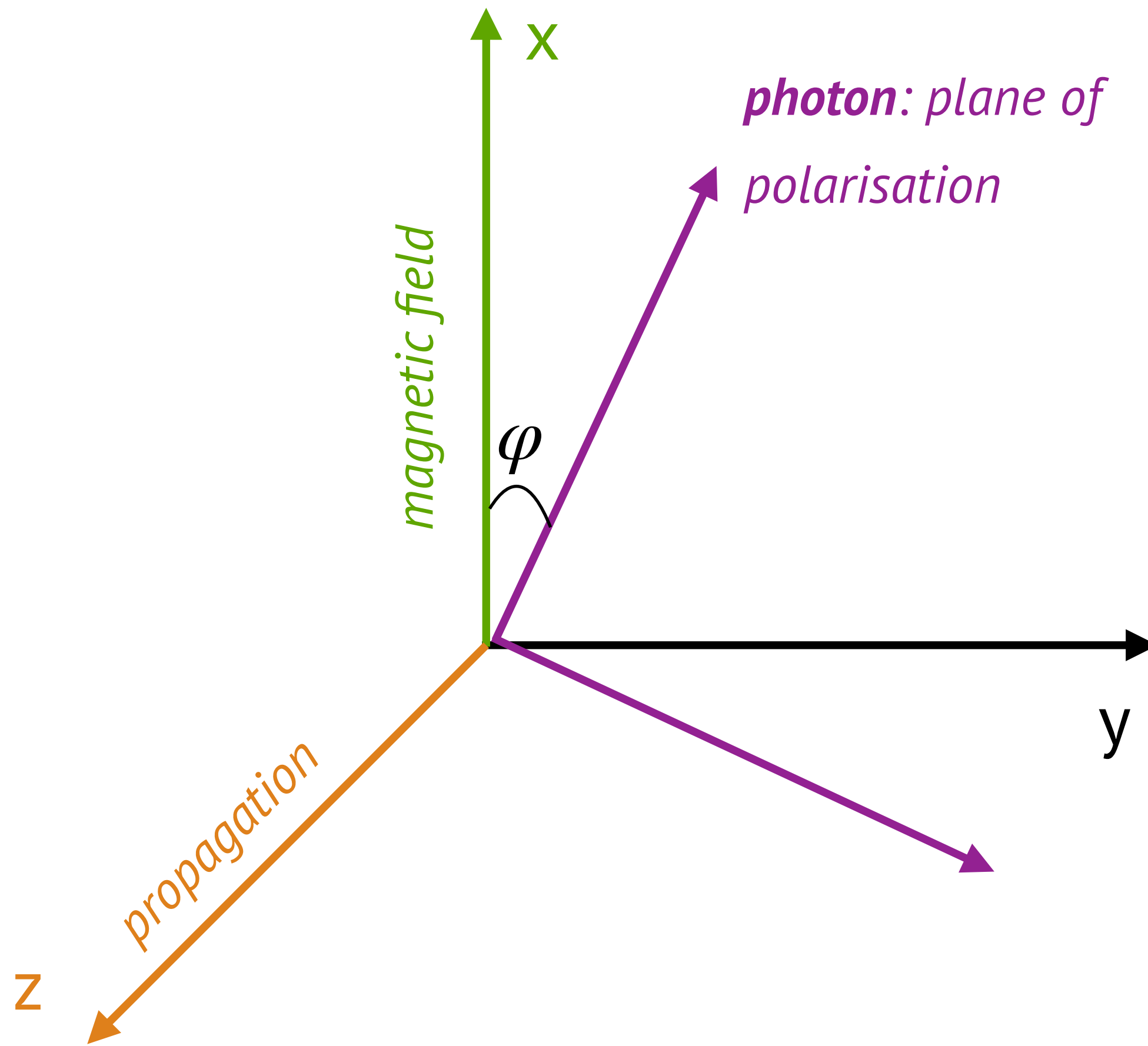
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source

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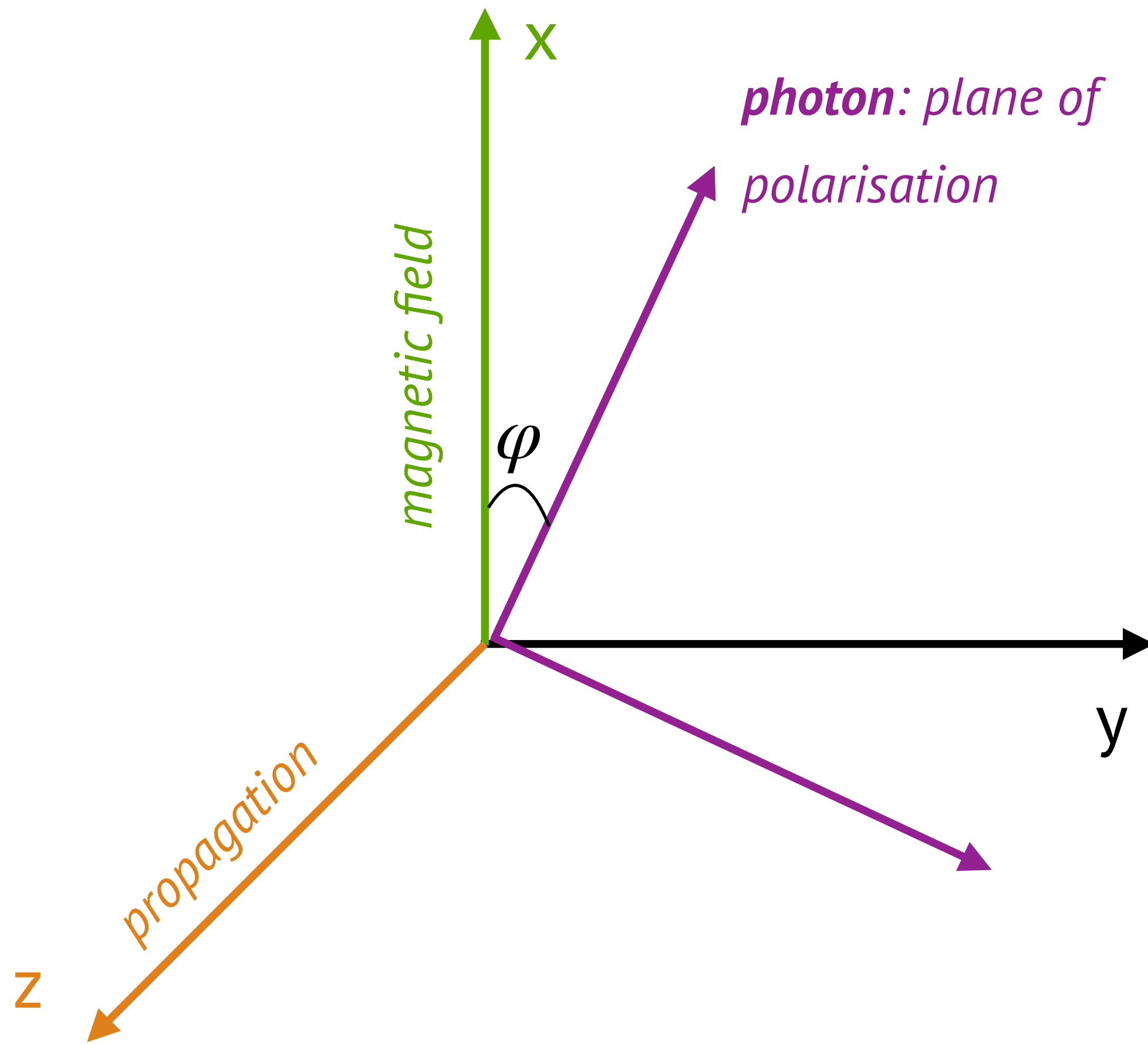
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three-state quantum system

$$\mathcal{A}(z) = \begin{pmatrix} A_1(z) \\ A_2(z) \\ a(z) \end{pmatrix} \begin{matrix} \rightarrow \text{photon} \\ \rightarrow \text{polarisations} \\ \rightarrow \text{ALP field} \end{matrix}$$

source

- ▶ initial ALP state needed to build wave function
- ▶ initial photon polarisation also required (A_1, A_2)
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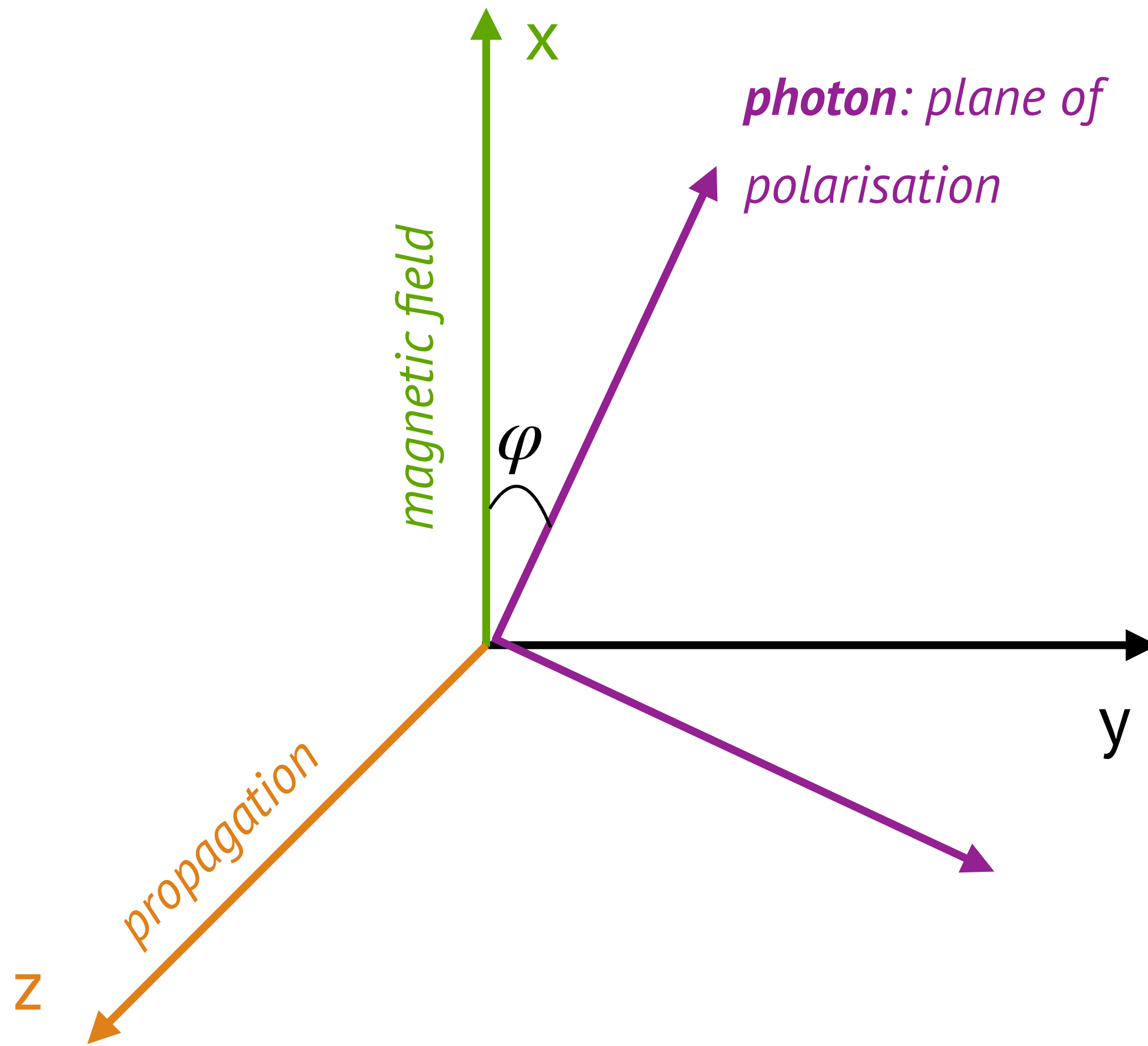
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```
source.add(alp.SourceALPState(0., 0.))  
source.add(alp.SourceNoPolarisation())
```



mixing matrix

$$\mathbb{M} = \begin{pmatrix} \Delta_{\parallel} \cos^2 \varphi + \Delta_{\perp} \sin^2 \varphi & (\Delta_{\parallel} - \Delta_{\perp}) \sin \varphi \cos \varphi & \Delta_{a\gamma} \cos \varphi \\ (\Delta_{\parallel} - \Delta_{\perp}) \sin \varphi \cos \varphi & \Delta_{\parallel} \sin^2 \varphi + \Delta_{\perp} \cos^2 \varphi & \Delta_{a\gamma} \sin \varphi \\ \Delta_{a\gamma} \cos \varphi & \Delta_{a\gamma} \sin \varphi & \Delta_a \end{pmatrix}$$

depends on

- $\Delta_{a\gamma}$ → B: external magnetic field
- $\Delta_{a\gamma}$ → $g_{a\gamma}$: coupling constant
- Δ_a → ma : ALP mass
- Δ_{\perp} → E: photon energy
- Δ_{\parallel} → ne : plasma density in medium

equation of motion

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density = alp.PlasmaDensityUniform(densityIGM)
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ALP-photon mixing

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ALP-photon mixing

```
alpPhoton = alp.ALPhotonMixing(axionMass * eV / c_light ** 2, coupling / eV, bField, density, .5)
```


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- ▶ random numbers compared to probability of photon-ALP conversion to decide if oscillation occurs

For a homogeneous magnetic field and medium, the equations of motion can be solved analytically to determine the conversion probability.

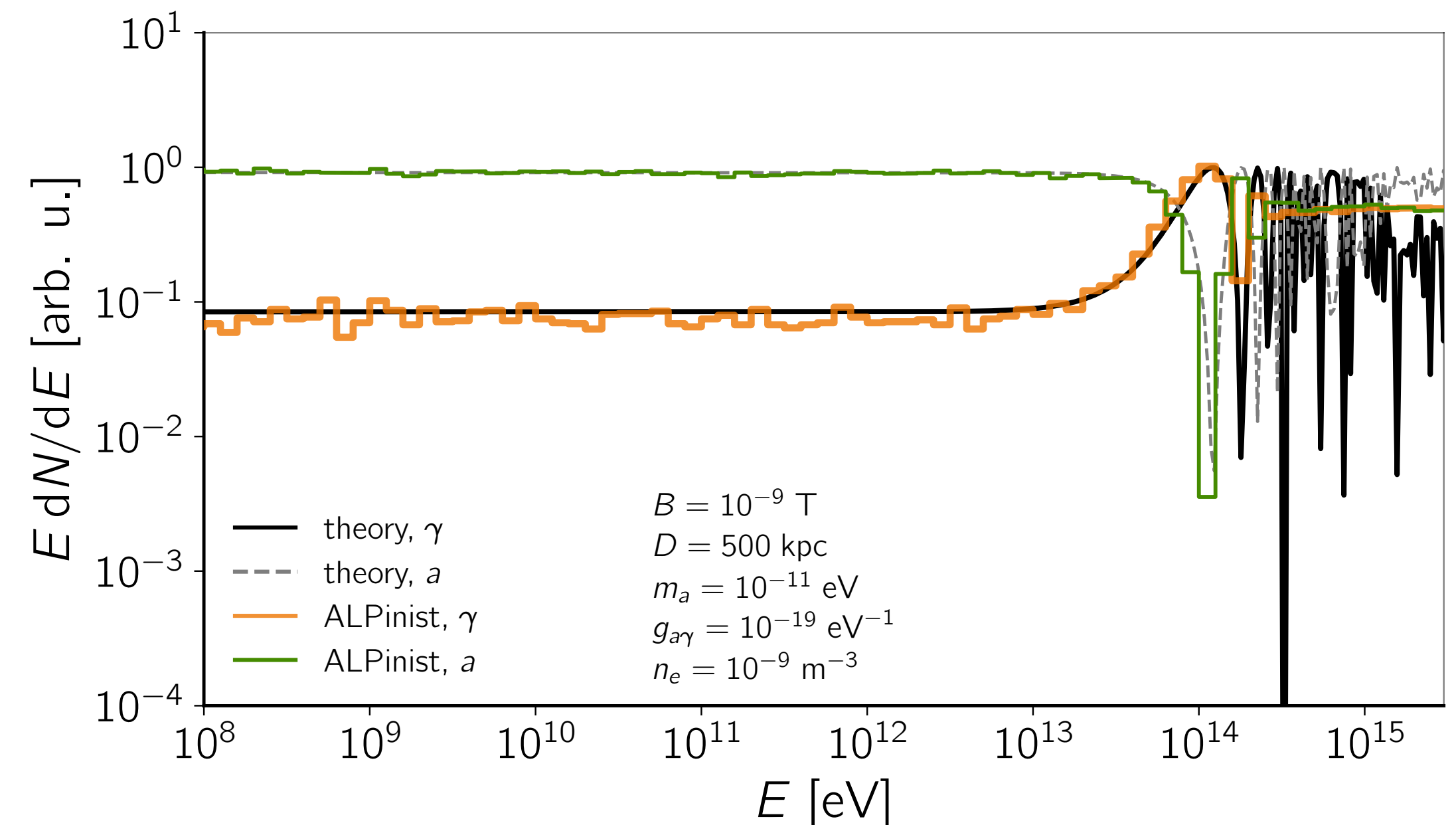
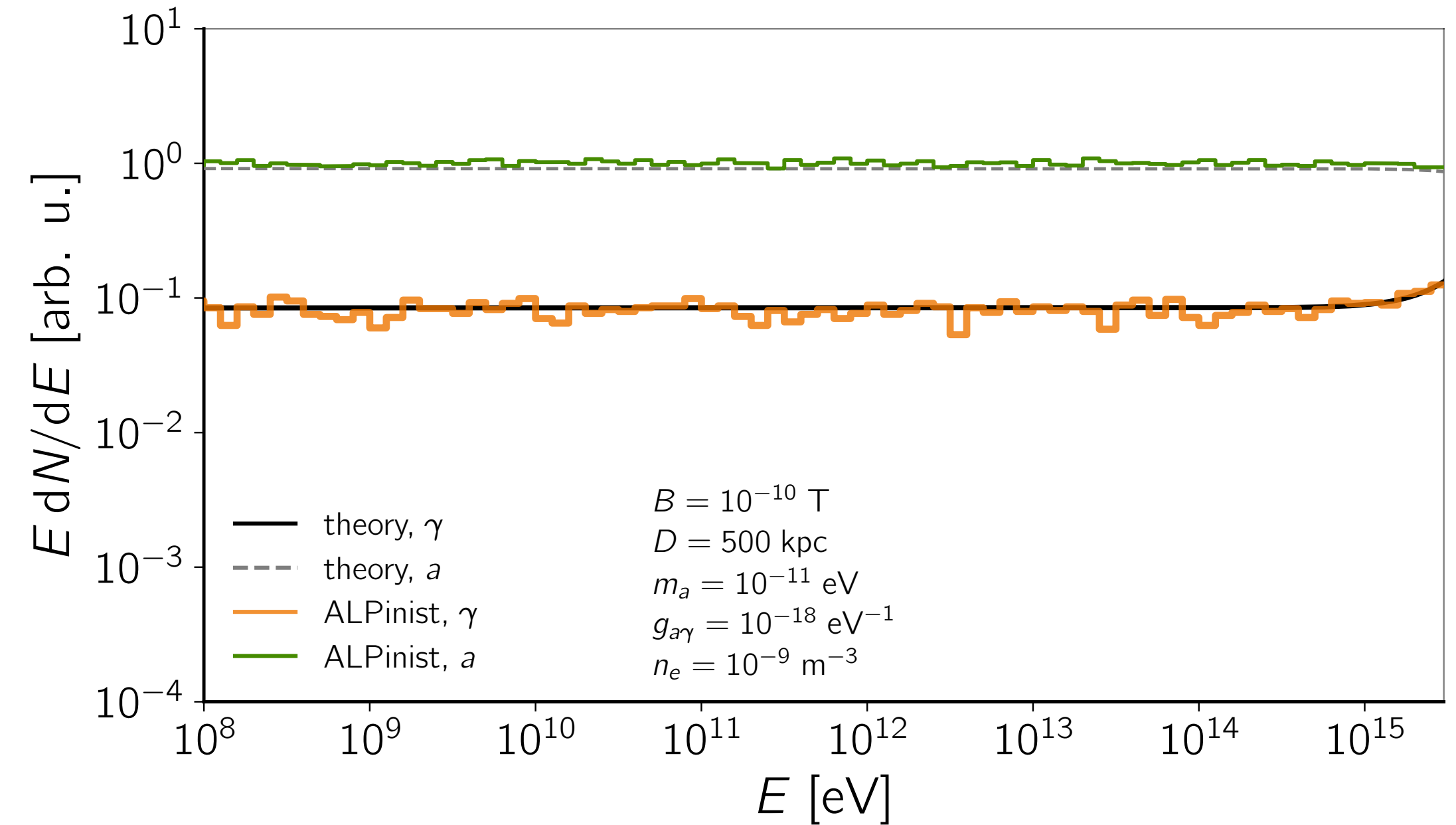
conversion probability

$$P_{a\gamma}(z) = \sin^2 \theta \sin^2 \left(\frac{z}{2} \sqrt{(\Delta_a - \Delta_{pl})^2 + 4\Delta_{a\gamma}} \right)$$

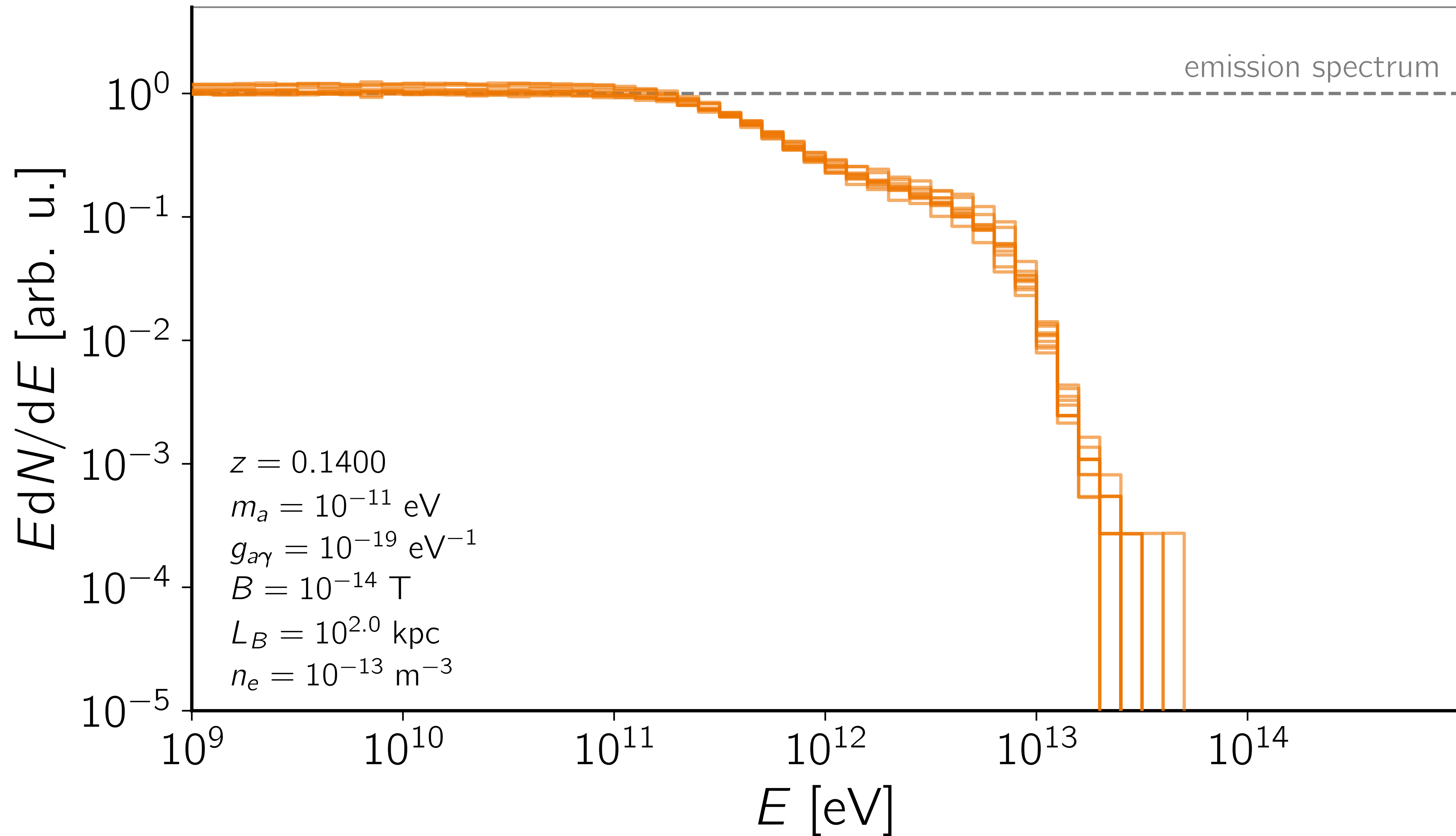
mixing angle

$$\theta = \frac{1}{2} \arcsin \left(\frac{2\Delta_{a\gamma}}{\sqrt{(\Delta_a - \Delta_{pl})^2 + 4\Delta_{a\gamma}}} \right)$$

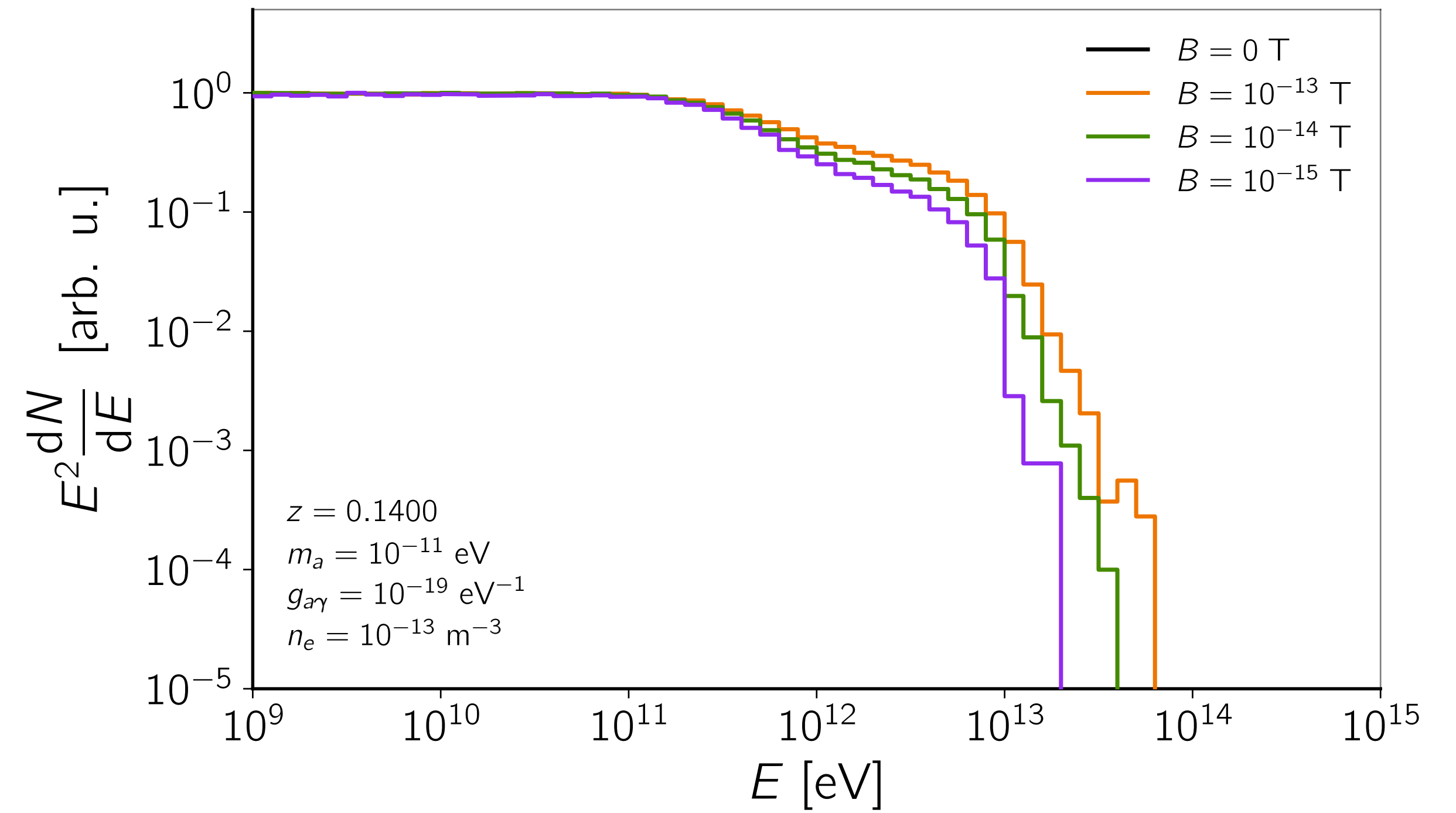
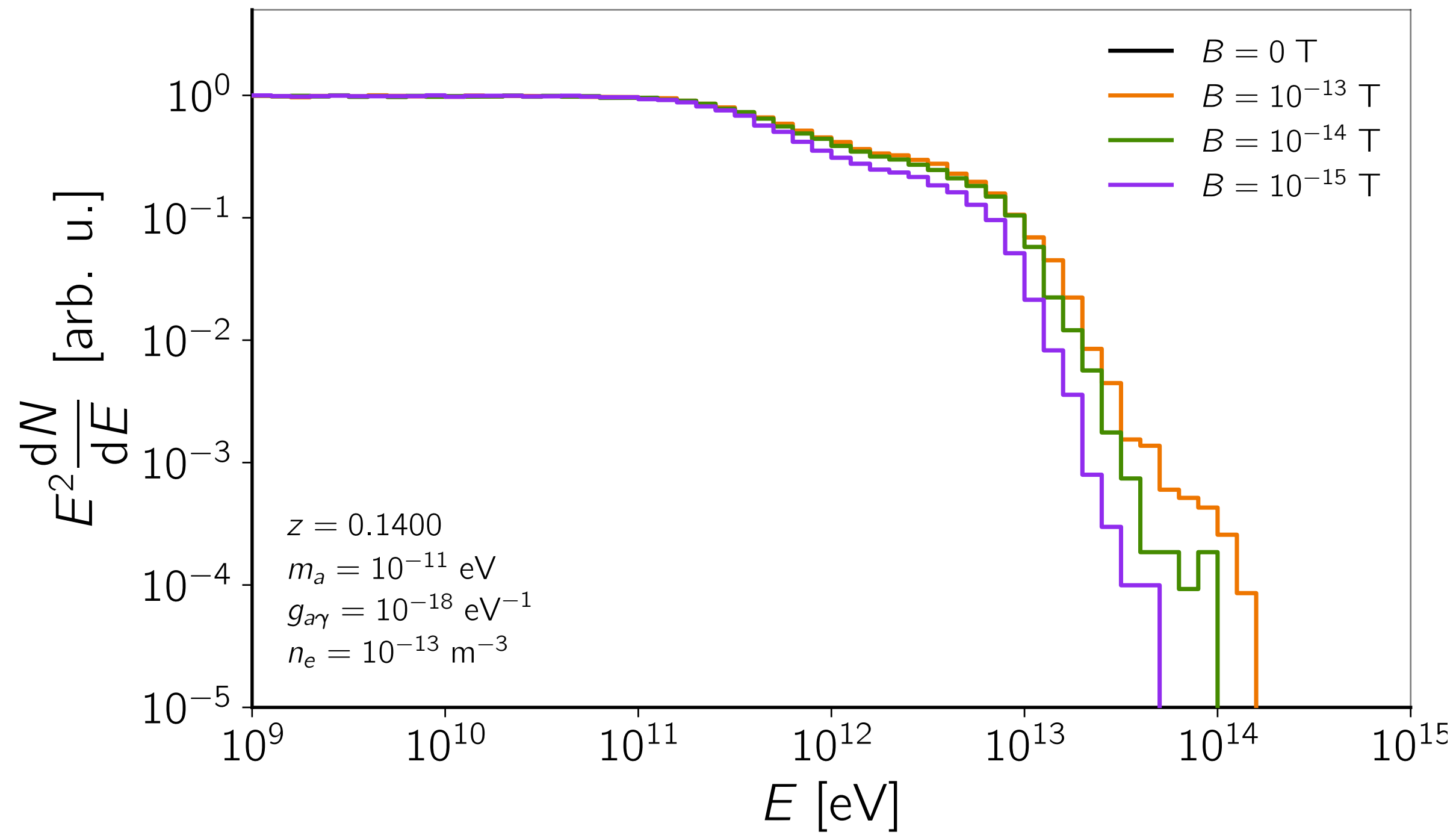
For $\theta=45^\circ$ mixing is maximal and an oscillatory behaviour sets in.



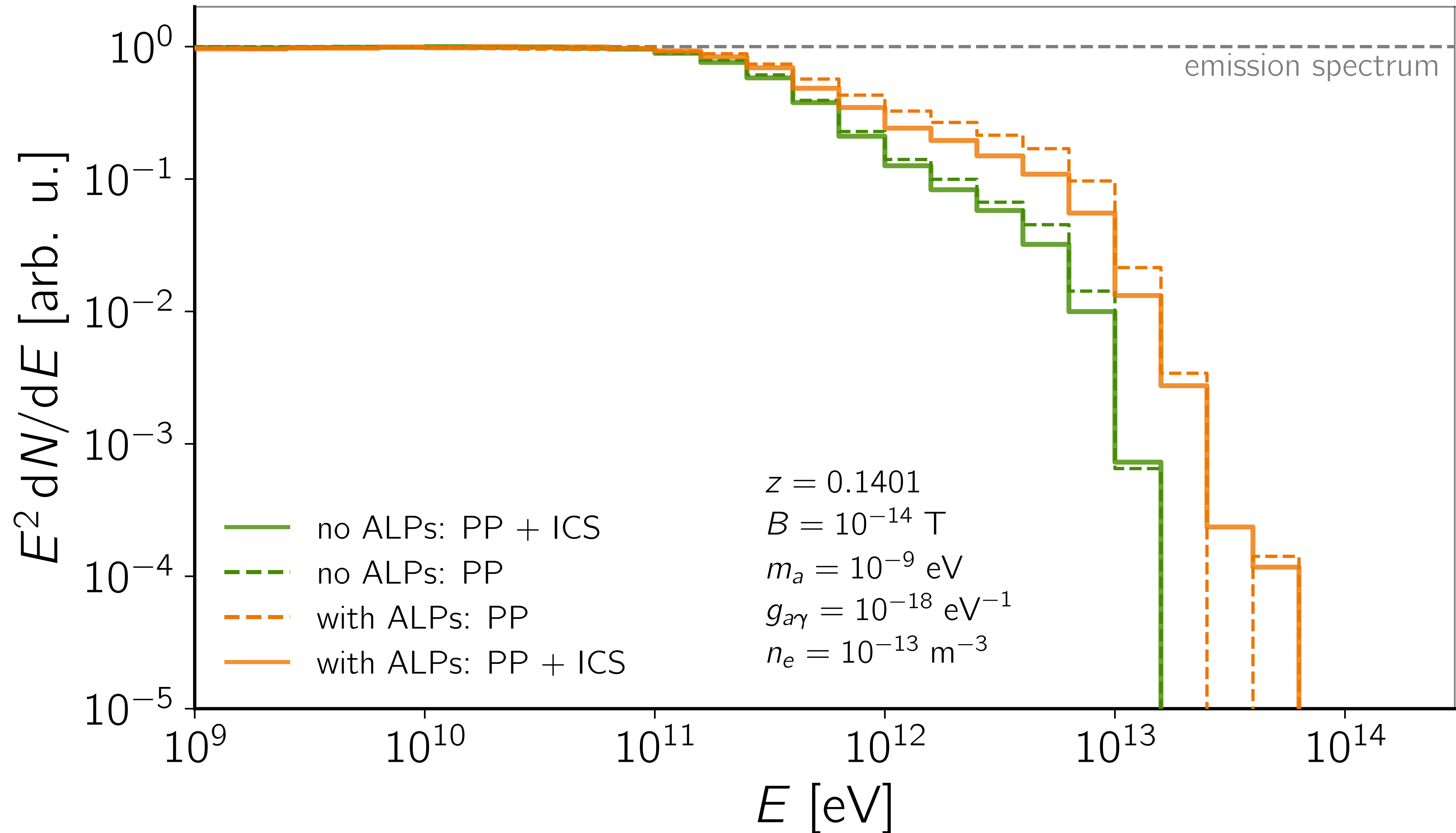
simulations with turbulent magnetic fields



simulations with pair production and turbulent magnetic fields



adding inverse Compton scattering



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- ▶ first-ever simulations of electromagnetic cascades including all effects: pair production + inverse Compton scattering + ALP-photon mixing
 - ◆ cascades initiated via inverse Compton can provide an additional contribution to the gamma-ray flux